

# *Dynamics of Transient Liquid Injection:*

*K-H instability, vorticity dynamics, R-T instability,  
capillary action, and cavitation*

**William A. Sirignano**

**University of California, Irvine**

- **Round liquid columns without aerodynamic effects**
- **Stretched “conical” and fan jets without aerodynamic effects**
- **Planar free films with aerodynamic effects**
- **Round jets without surface tension and density discontinuity**
- **Round jets with surface tension and density discontinuity**
- **Liquid segment analysis**
- **Cavitation and bubble growth, distortion, and collapse**

*Support from US Army Research Office*

*Collaborations: former students, Dr C. Mehring and Dr. S. Dabiri; current student, Ms. D. Jarrahbashi; Prof. R. H. Rangel, and Prof. D. D. Joseph.*

# Stability of Cylindrical Liquid Column: Rayleigh Analysis

Surface area and surface energy of resulting drops are 79% of original cylinder values. Internal liquid pressure increases by about 6% causing a few percent increase in liquid enthalpy. Remaining enthalpy increase occurs due to viscous dissipation.

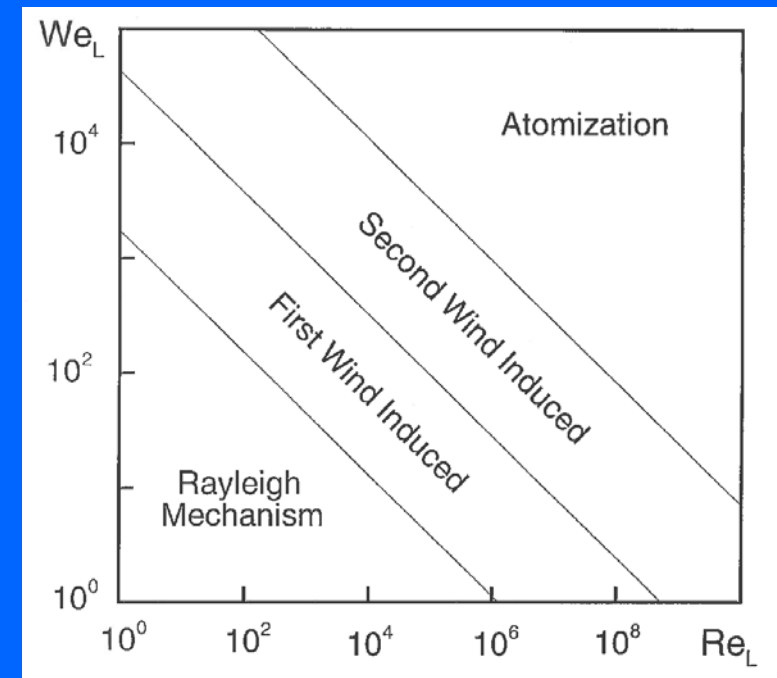
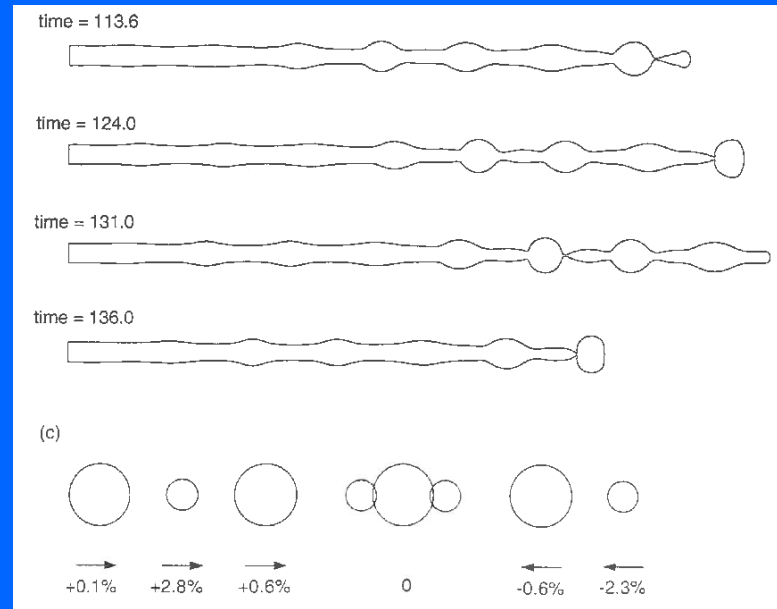
$$\frac{S_{drop}}{S_0} = 2 \left( \frac{R_{drop}}{R_0} \right)^2 \left( \frac{\lambda}{R_0} \right)^{-1} = 0.7933$$

$$\frac{2R_0}{R_{drop}} = 2 \left( \frac{4 R_0}{3 \lambda} \right)^{1/3}$$

$$= 1.058$$

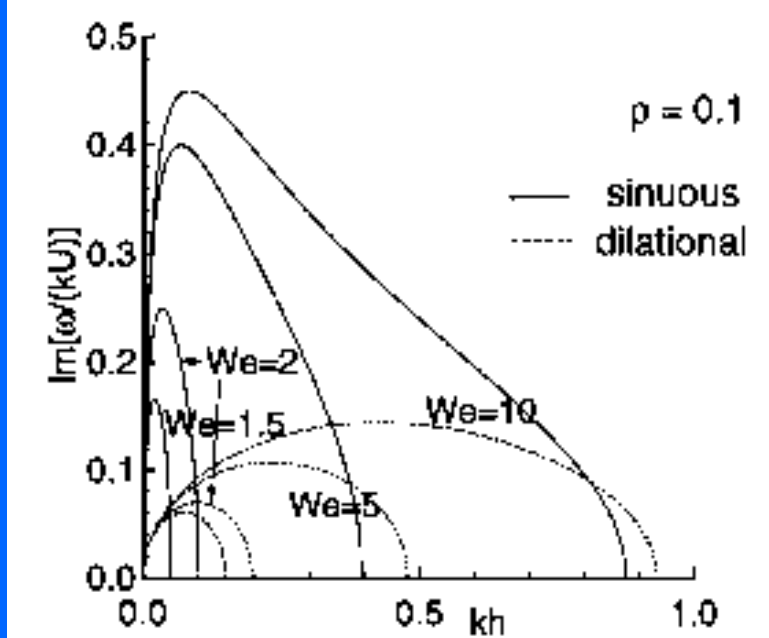
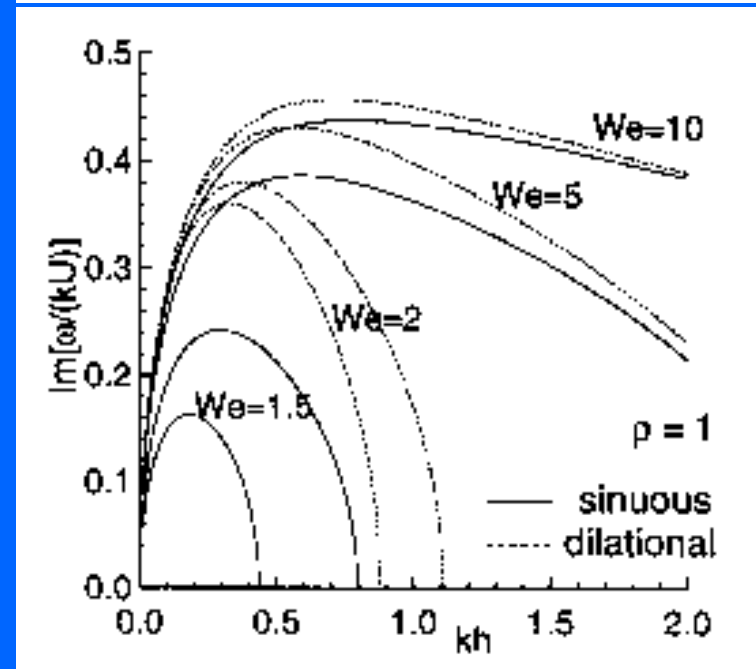
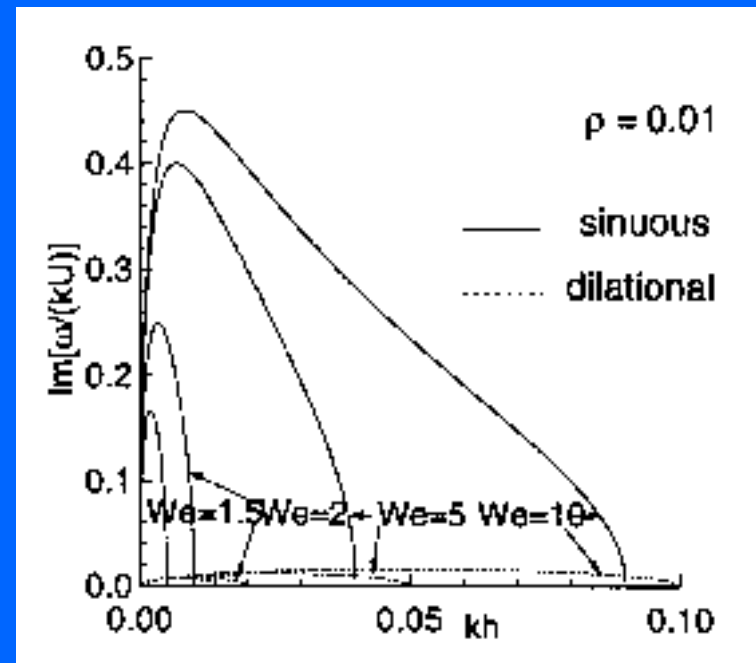
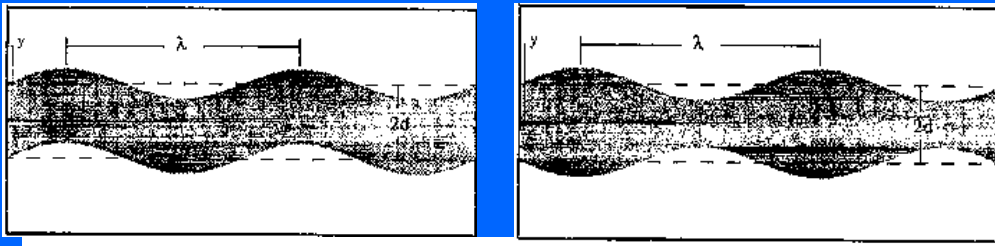
A more accurate nonlinear analysis predicts small satellite droplets also occur.

To obtain many small droplets (i.e., spray) and increase surface energy, kinetic energy of surrounding stream must be exchanged.



# Capillary / Kelvin-Helmholtz instability of planar liquid sheets

Linear theory predicts the dilational mode becomes more probable than the sinuous mode at high pressures



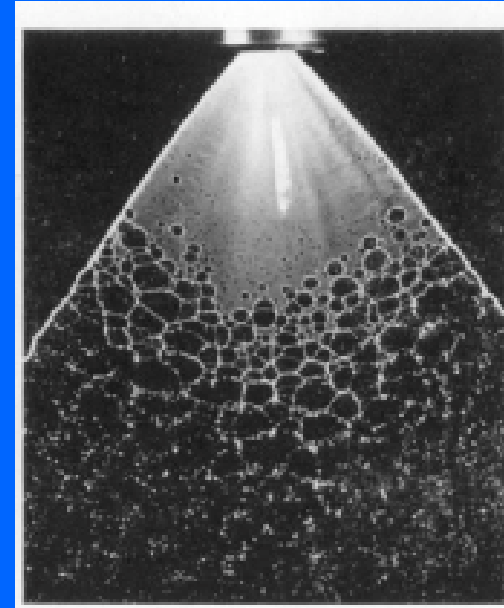
# *Capillary Action on Thin Sheets: e.g., fan jet and “conical” jet*

**Dilational oscillations of thin free liquid films or sheets can result in local tearing.**

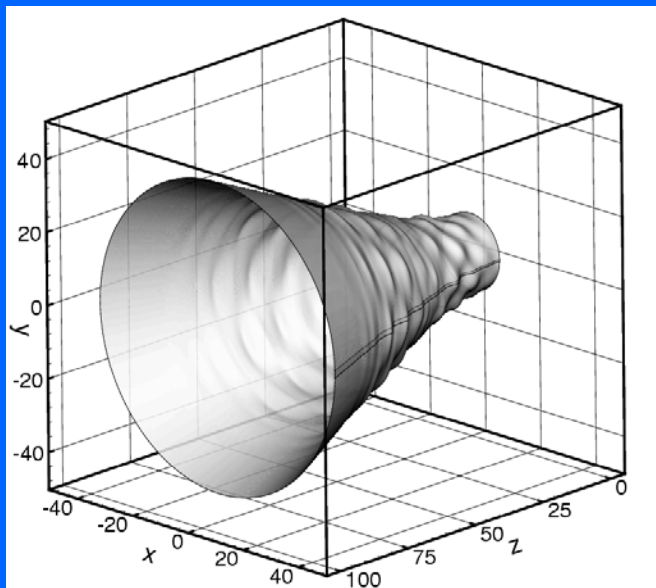
**No aerodynamic effects here.**

**Linear theory predicts dilational oscillations are more likely at high pressure.**

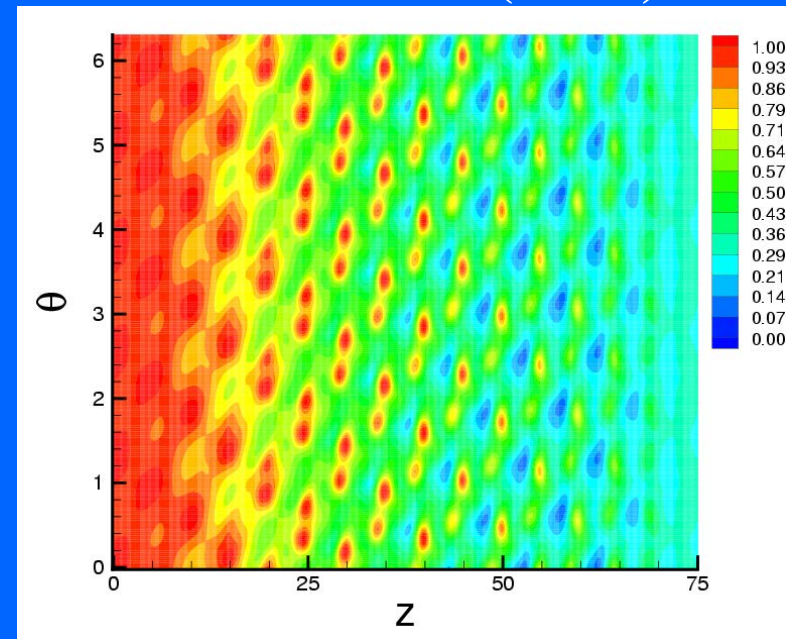
**Tearing of thin lobes formed in secondary instabilities will be addressed later.**



**Dombrowski & Fraser (1954)**

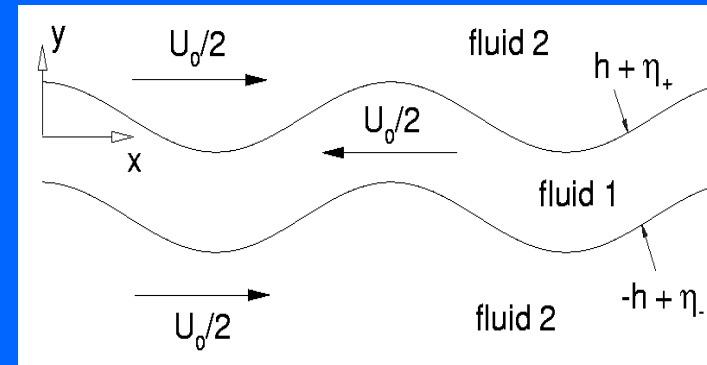


**Mehring & Sirignano (2004)**



## *Vortex Dynamics at Liquid-Gas Interface in Planar Flow*

The phase interface is a vortex sheet with velocity discontinuity. Examine nonlinear Kelvin-Helmholtz instability with surface tension and density jump.



Approximate sheet as a packed linear array of “vortex blobs.”

$$\frac{d(\Delta\Gamma)}{dt} = 2A \left( \frac{d\mathbf{u}}{dt} \right)_s \Delta s + \frac{A}{4} \frac{\partial}{\partial s} \left( \frac{\Delta\Gamma}{\Delta s} \right)^2 \Delta s - \frac{2}{We} \frac{\partial \kappa}{\partial s} \Delta s,$$

$$A = (\rho_2 - \rho_1) / (\rho_2 + \rho_1); \quad We = (\rho_2 + \rho_1) \lambda (\Delta U)^2 / \sigma$$

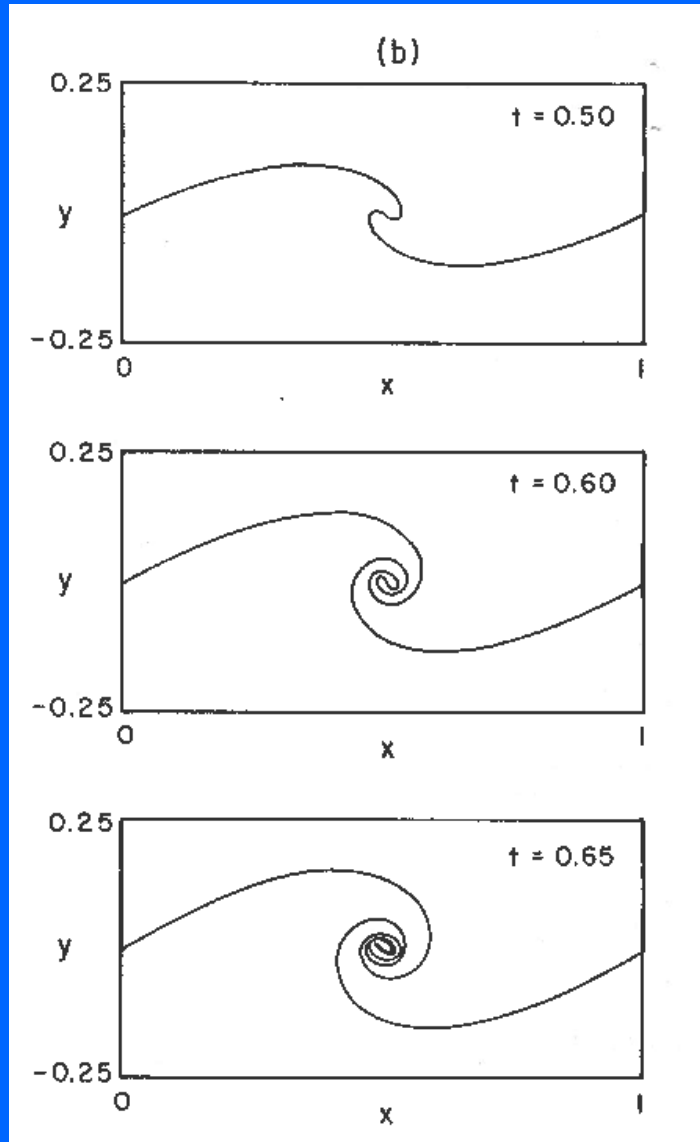
Zalosh (1976) and Rottman & Olfe (1977)

Circulation changes with time due to surface tension.

Also, density difference causes a temporal change in circulation.

Similar qualitative behavior might be expected for axisymmetric flow.

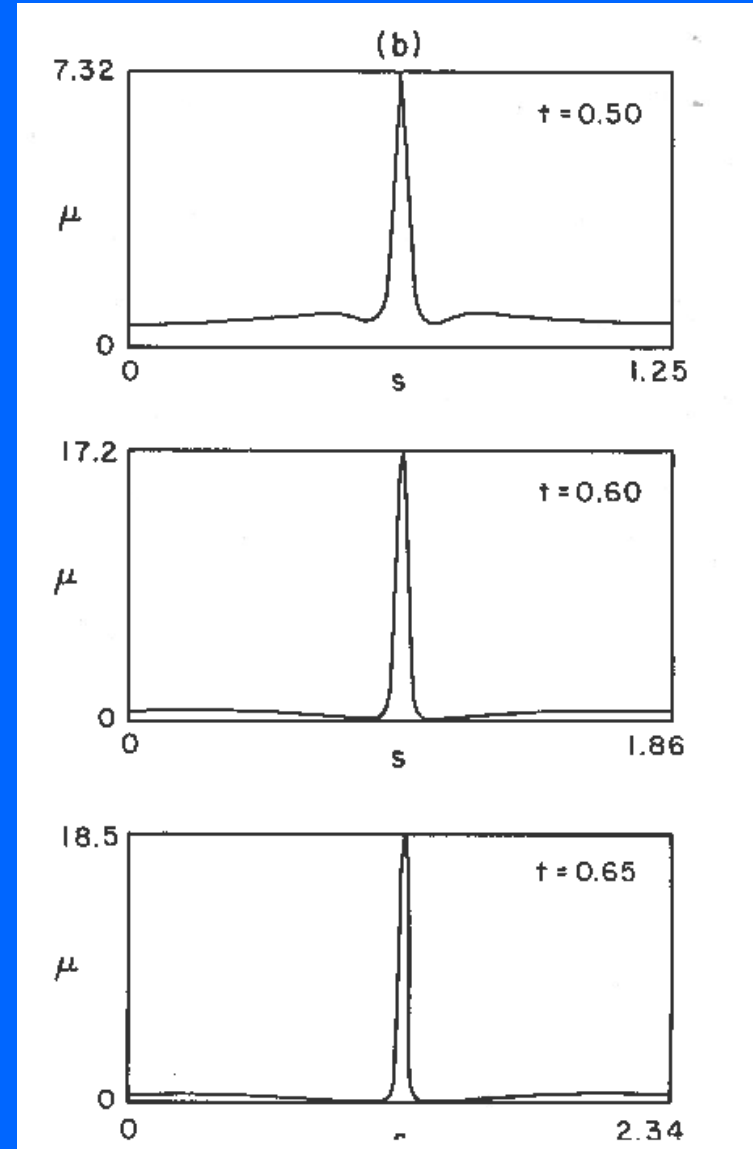
# *Surface contour and vorticity with no surface tension and no density difference*



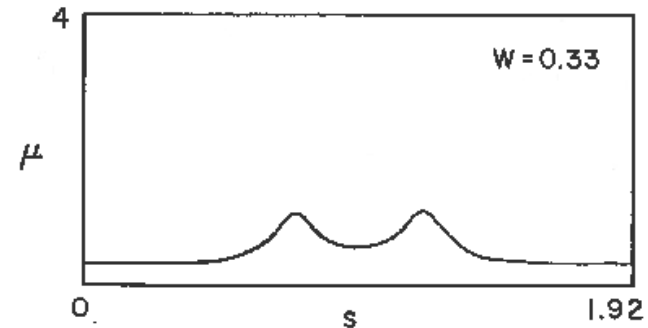
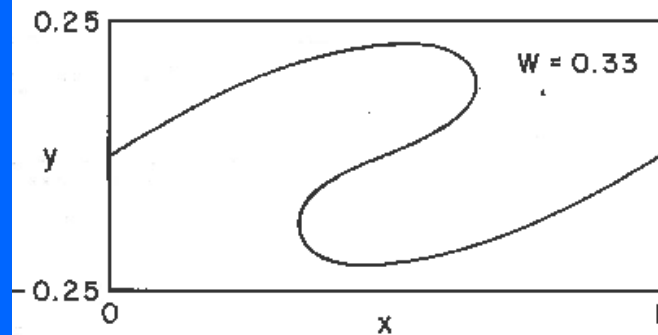
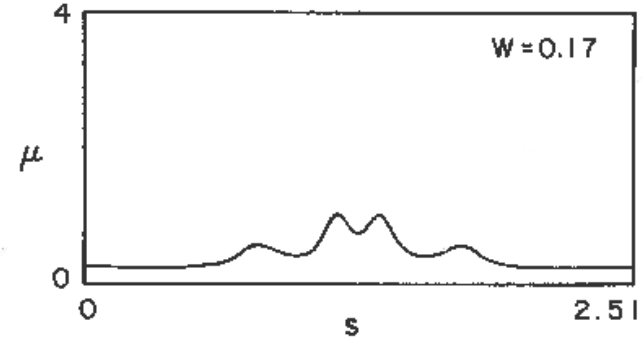
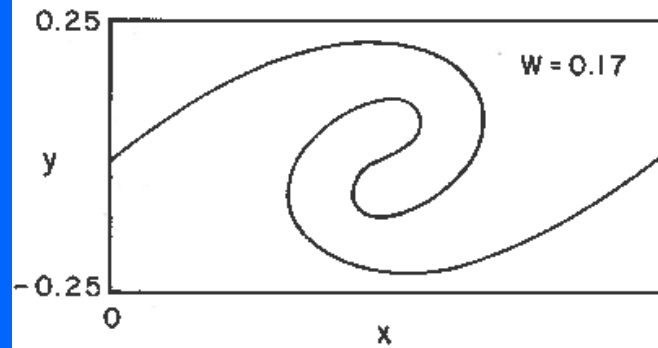
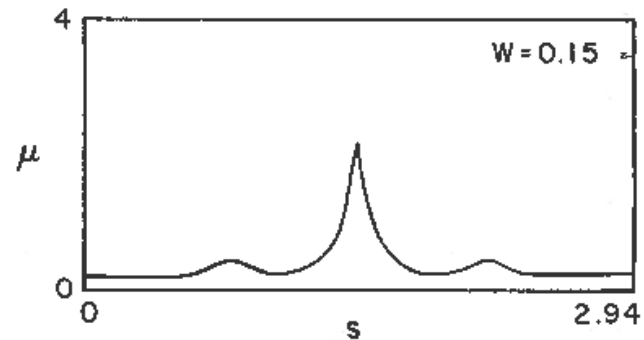
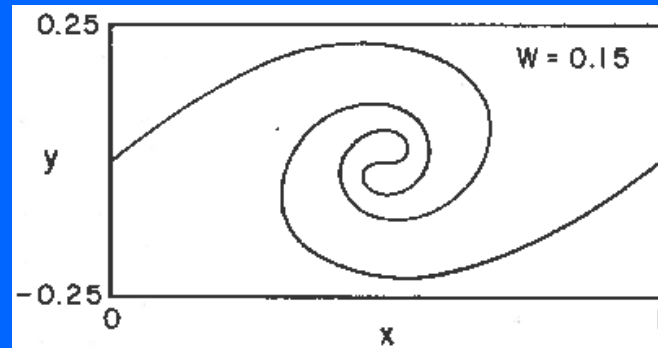
← Contour

Vorticity →

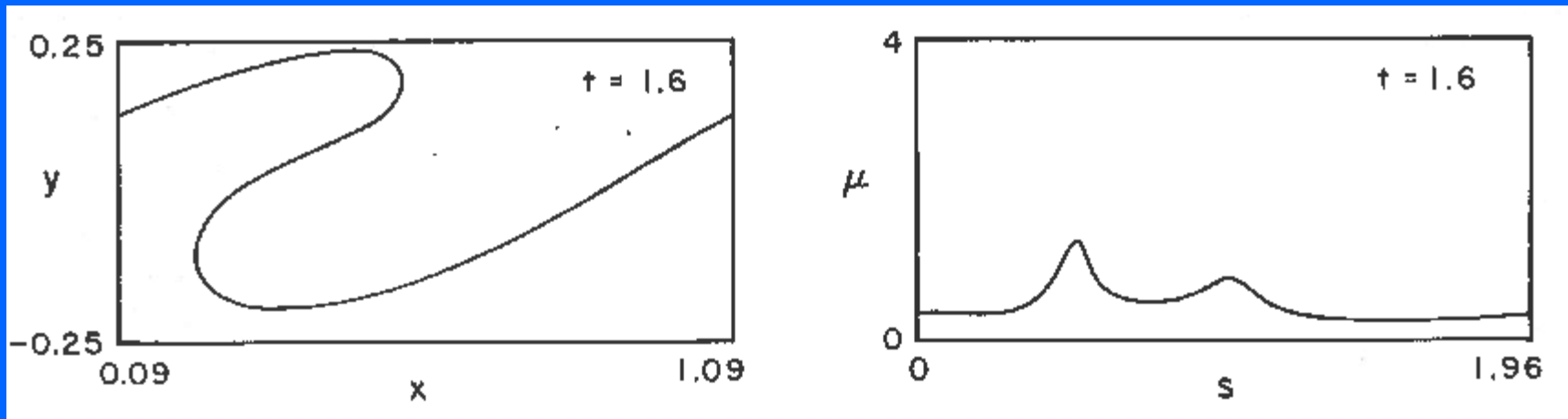
**Rangel &  
Sirignano  
(1988)**



# *Surface tension inhibits roll-up and vorticity concentration: no density difference*



*Density reduction also inhibits roll-up,  
here in combination with surface tension*



- “Cone” shape surface results from surface tension and density discontinuity.
- Similar shape distortion from KH instability with surface tension and density discontinuity is seen for axisymmetric case, i.e., round jet.
- 3D distortions on these cones can be expected.



*Planar free film with K-H instability and capillary action. A cone shape can develop in time with concentration of vorticity.*

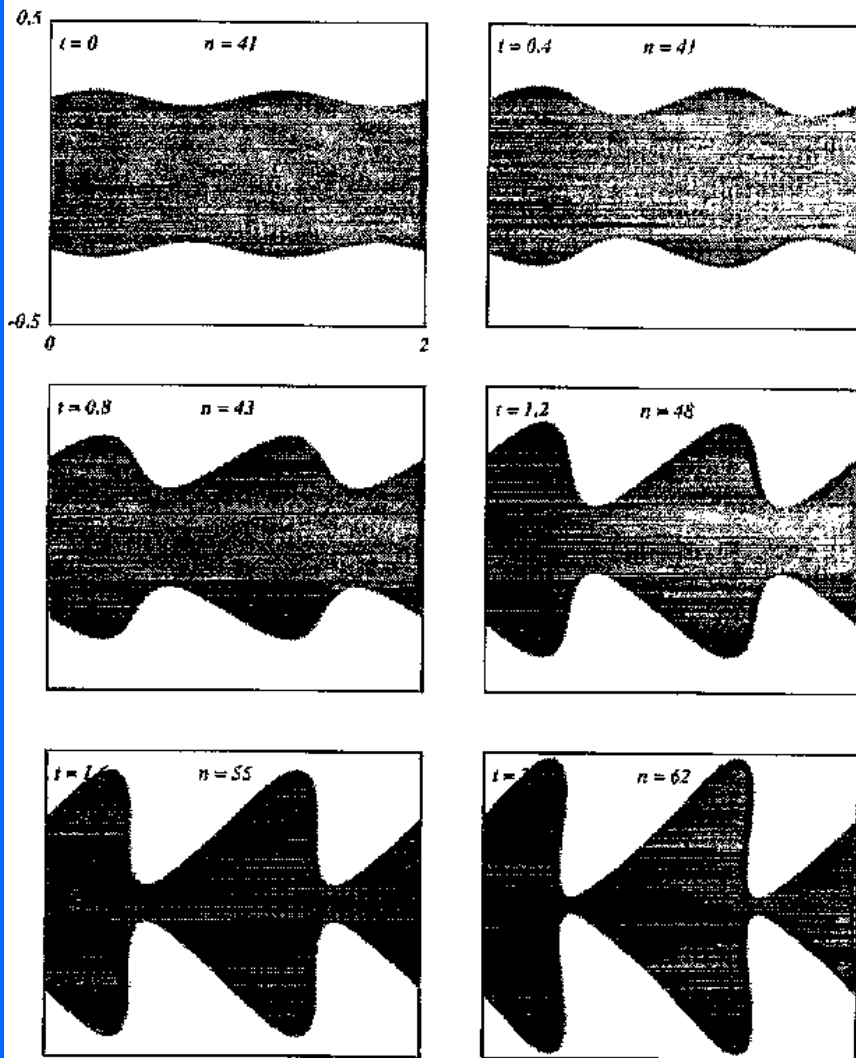


FIG. 14. Time evolution of the dilational mode for  $\rho = 1$ ,  $W = 0.67$ , and  $h = 0.25$  ( $H = 0.373$ ).

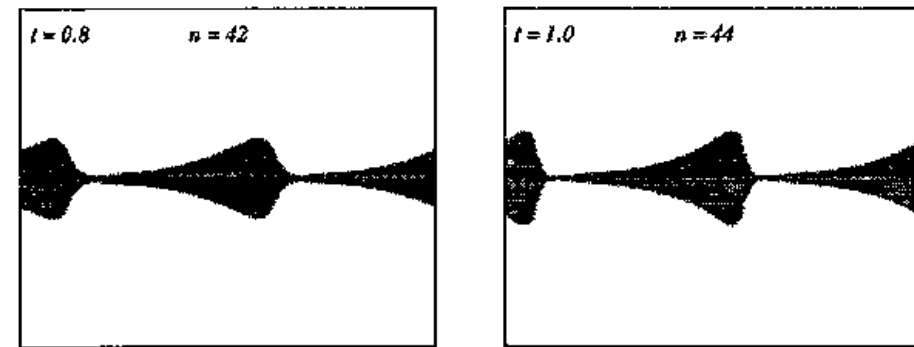


FIG. 8. Time evolution of the dilational mode for  $\rho = 0.25$ ,  $W = 0.5$ , and  $h = 0.05$  ( $H = 0.1$ ).

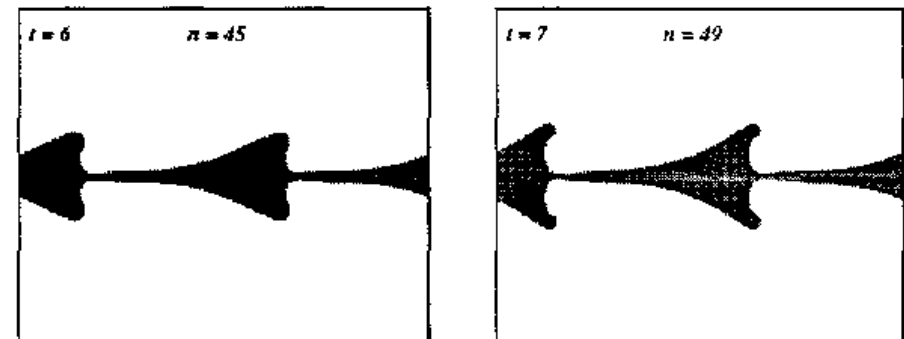
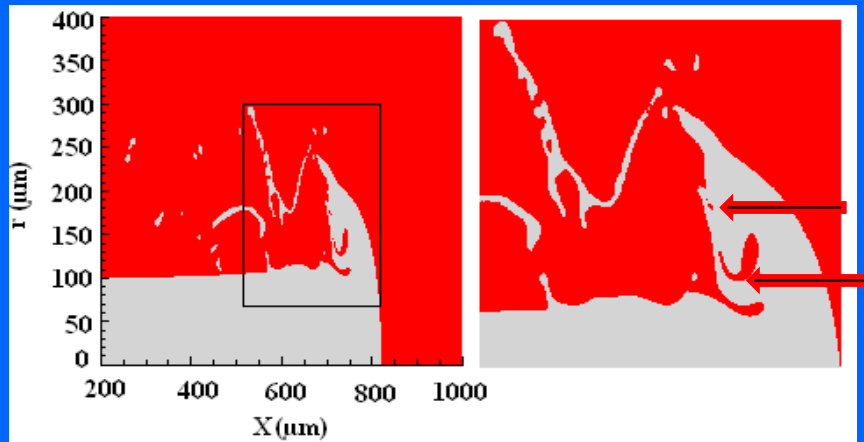


FIG. 11. Time evolution of the dilational mode for  $\rho = 0.25$ ,  $W = 0.05$ , and  $h = 0.05$  ( $H = 0.1$ ).

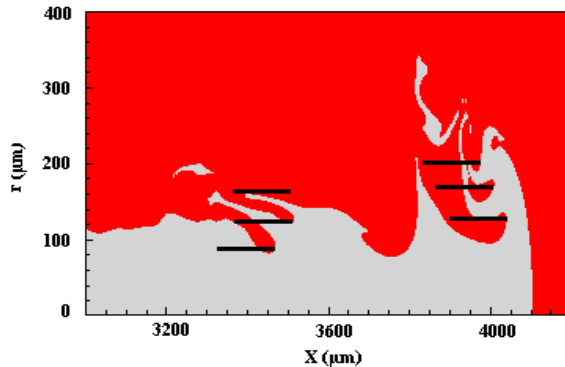
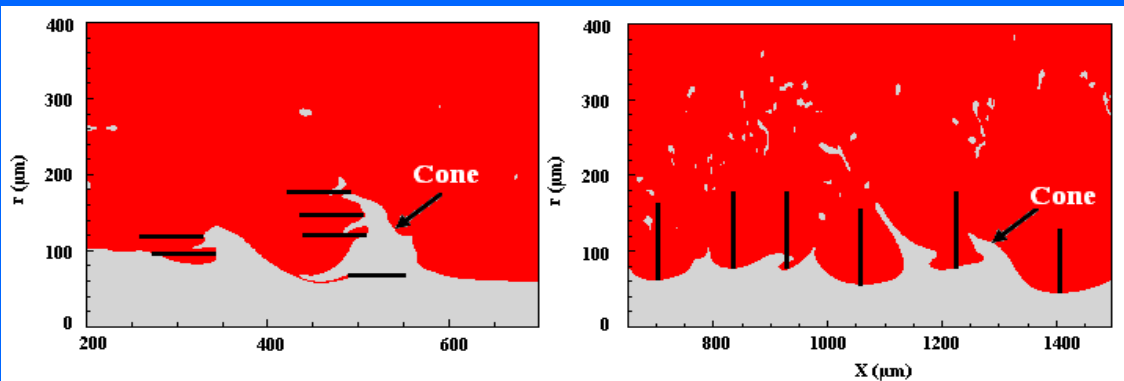
**Rangel & Sirignano (1991)**

# Axisymmetric round jet with K-H and R-T instabilities



$Re=16,000$ ,  $We=230,000$ ,  $t=40 \mu s$

- Start-up: Heavier fluid on cap accelerates into the lighter fluid; R-T instability on the rear side of the jet cap.
- R-T instability appear on top of the K-H cones drawn out of the liquid.



$Re=16,000$ ,  
 $We=230,000$ ,  
 $t=100 \mu s$

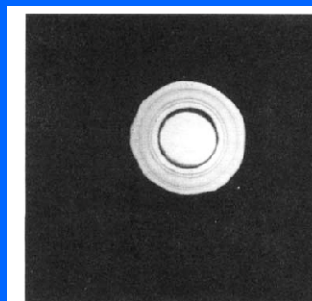
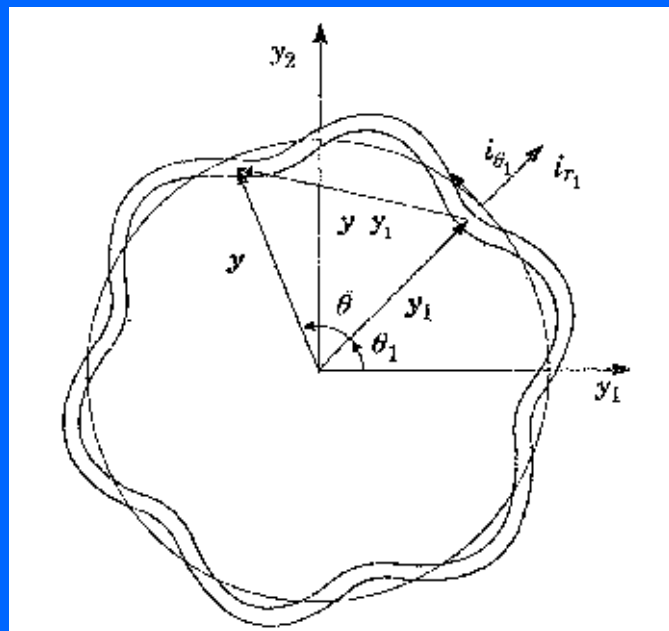
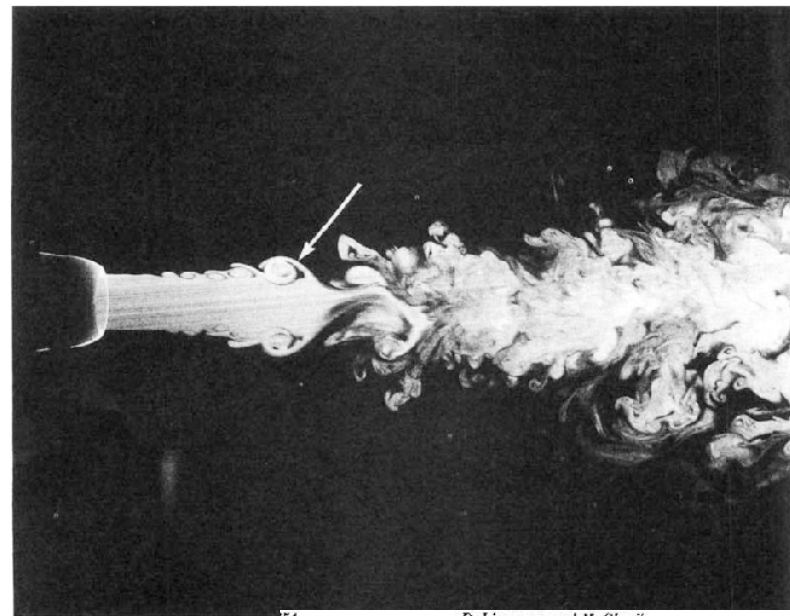
Jarrahbashi &  
Sirignano (2013)

# Round jets and unstable vortex rings

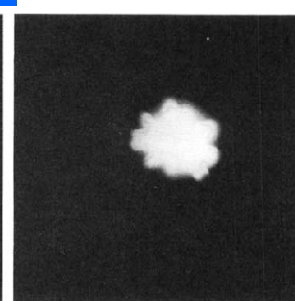
Axisymmetric vortex rings develop  
3D instabilities

Liepmann & Gharib (1992),  
water-into-water

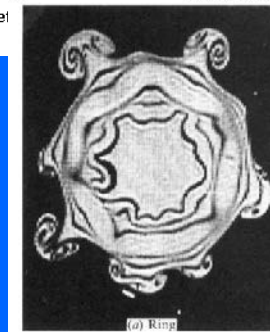
Widnall & Sullivan (1973),  
air-into-air



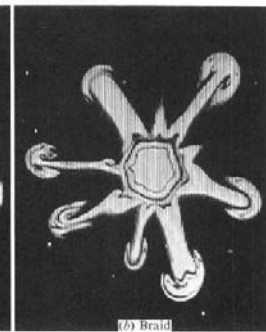
Ring



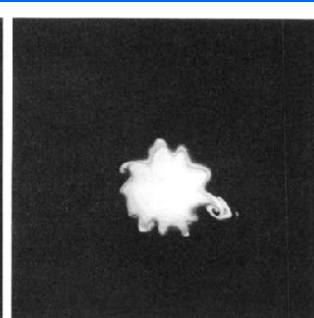
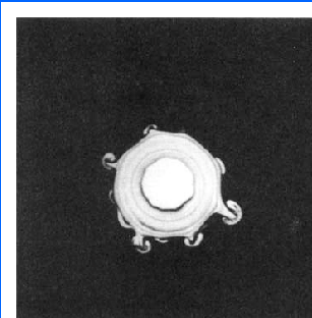
Braid

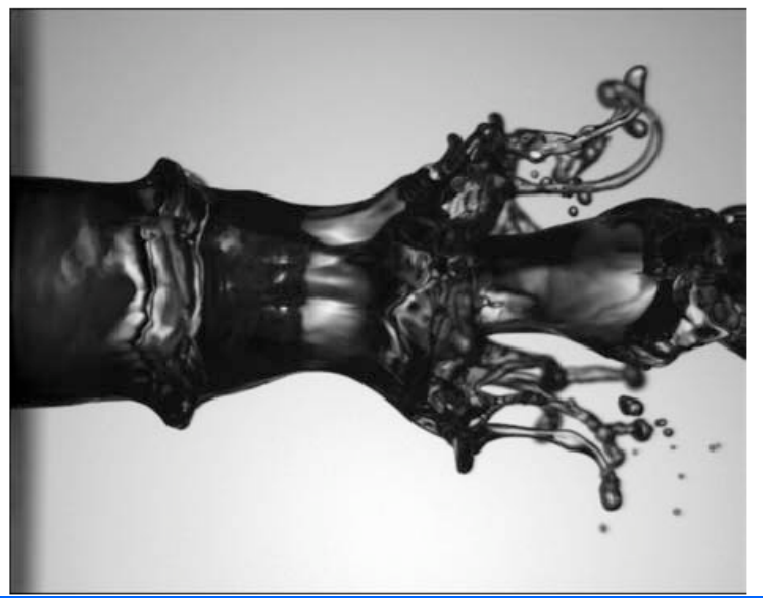


Ring



Braid





# *Liquid jet with co-axial air*

- Cones form.
- Azimuthal lobes form on cones with attribution to Rayleigh-Taylor instability.
- Ligaments form from lobes.

**Marmottant & Villermaux (2004)**

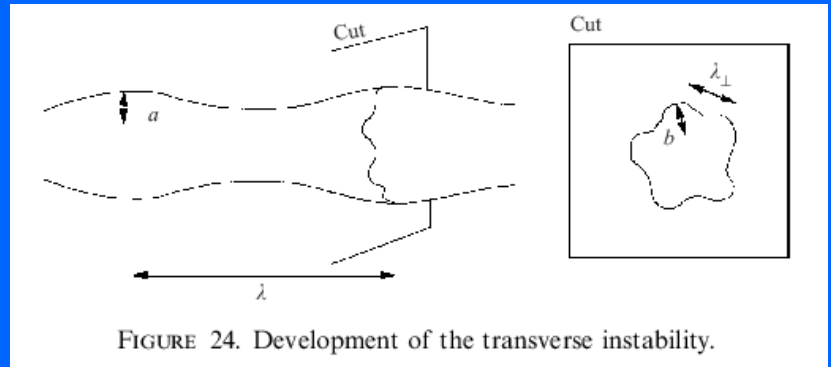
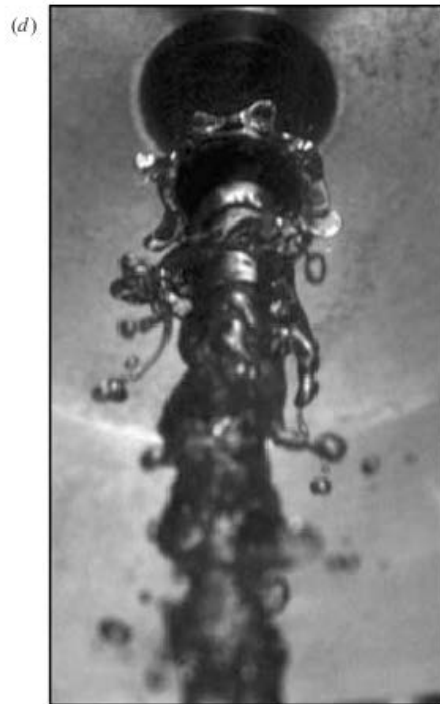
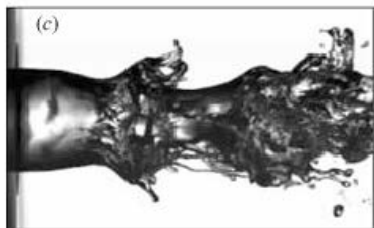
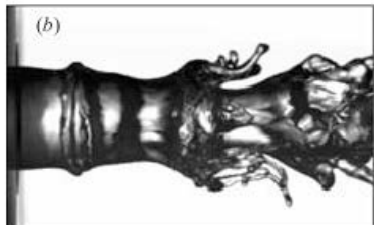


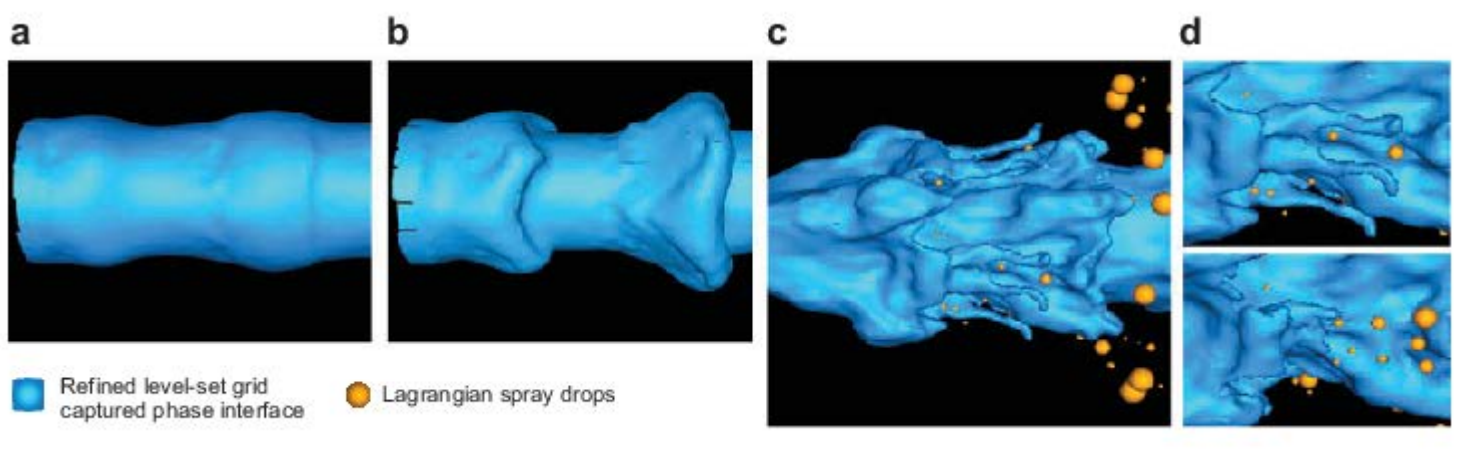
FIGURE 24. Development of the transverse instability.

$$\omega_{iRT} = \left( \frac{2}{3\sqrt{3}} \right)^{1/2} \left( \frac{\rho_1 g^3}{\sigma} \right)^{1/4}$$

$$\lambda_{RT} = 2\pi \left( \frac{3\sigma}{\rho_1 g} \right)^{1/2}$$

*How does vorticity dynamics act in lobe formation?*

# 3D Navier-Stokes / Level-Set Solutions



$$Re_g = 3770$$

$$We_g = 34$$

$$Re_l = 295$$

$$We_l = 0.6$$

Kim, Desjardins, Herrmann, and Moin (2007) – liquid jet with co-axial air flow, simulated M & V experiment.

- K-H axisymmetric waves deform to cone shape;
- Azimuthal distortion (attributed to R-T instability) yields lobes, then ligaments;
- Ligaments pinch-off via capillary action to yield droplets.

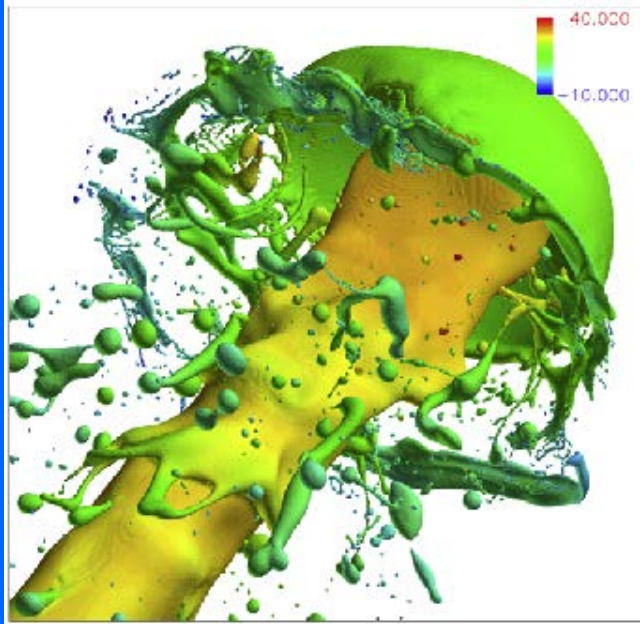
*Does 3D vorticity dynamics affect lobe formation? Are the lobes and ligaments “cousins” of the single-phase vortex-pairs?*

*Does R-T instability occur on ligaments also?*

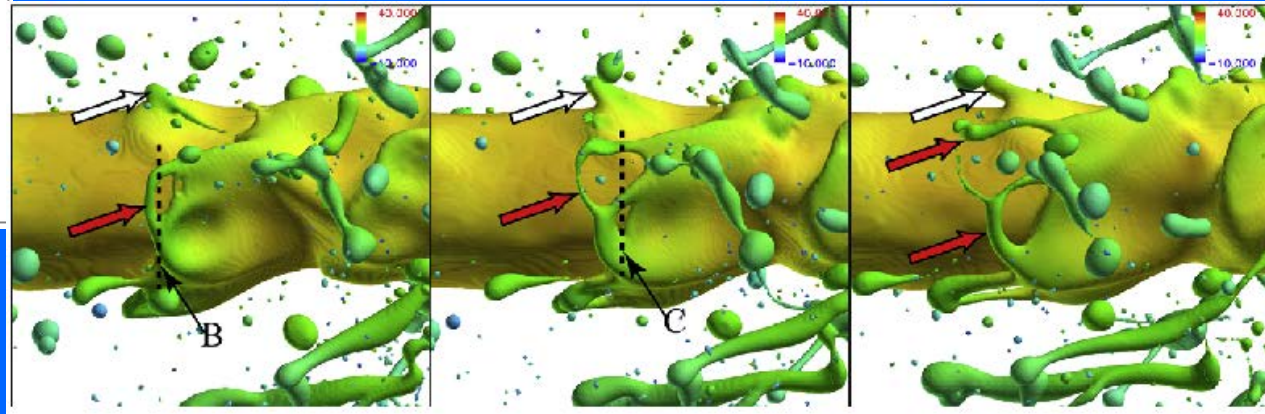


# *Liquid transient, round jet into still air*

*Shinjo & Umemura (2010), NS-LS  
 $Re = 440 - 1470$ ;  $We = 1270 - 14,100$*



- Cap forms on starting jet.
- Mass and size of cap increase.
- Instability causes shredding of cap.
- Lobes form on jet stem or core surface.

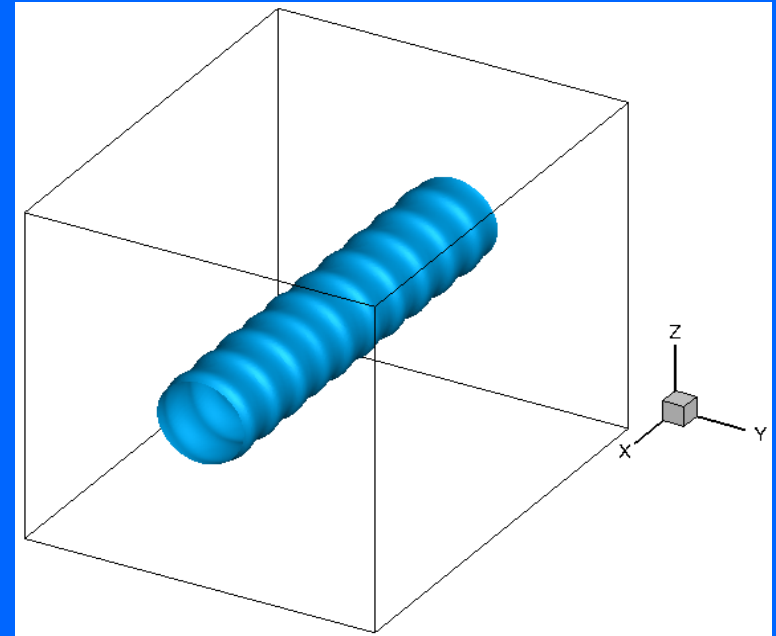


- Cones have azimuthal instability. Lobes form.
- Lobes stretch and tear under azimuthal instability.
- Ligaments form from torn rims.
- Capillary instabilities on ligaments. (*What about R-T instability on ligament?*)
- Droplets impact core surface.

# *Three-dimensional Liquid Jet Segment*

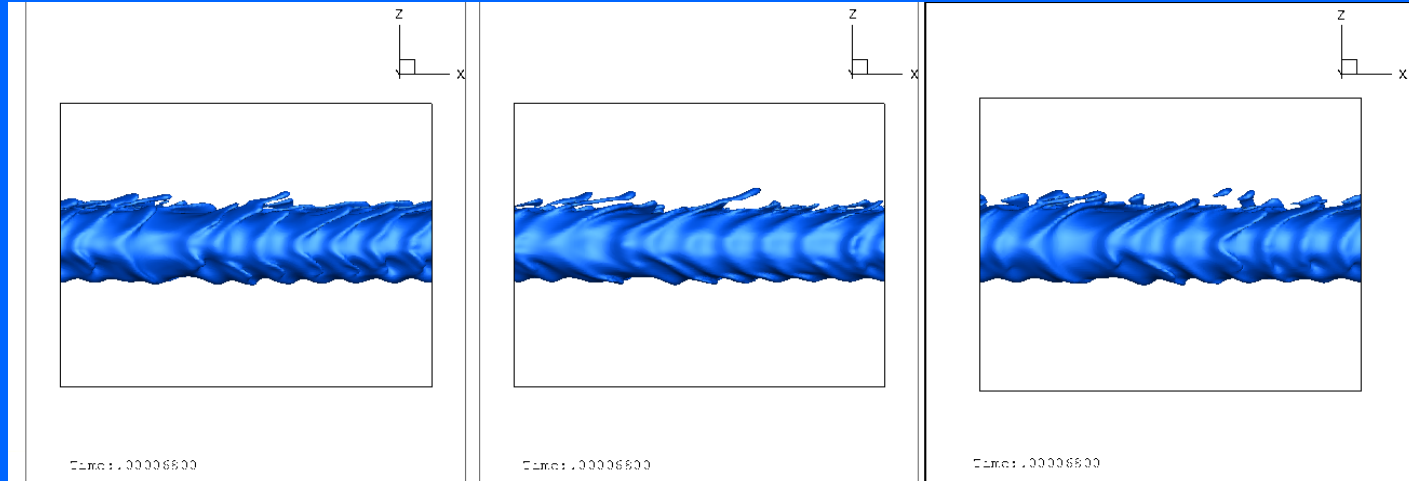
## *N-S / L-S solutions*

- The liquid-segment with 1 mm length **Jarrahbashi & Sirignano (2013)**
- and initial 200  $\mu\text{m}$  diameter.
- Continuity of the velocity and shear stress across the interface.
- Periodic boundary condition for all three components of velocity in the x-direction.
- Zero normal gradients of velocity and pressure on the outer  $xz$  and  $xy$  planes
- Initial disturbance is a sinusoidal surface wave with wavelengths 100  $\mu\text{m}$ .
- The initial wavelength was chosen to be consistent with full-jet results.
- Vortex rings form from the primary shear (K-H) instability at the jet interface.



# *Asymmetric 3D Distortion is not strongly dependent on initial conditions*

$t = 67 \mu s$



$Re = 1600$

$We = 230,000$

$\rho_1 / \rho_2 = 0.1$

$\mu_1 / \mu_2 = 9 \times 10^{-4}$

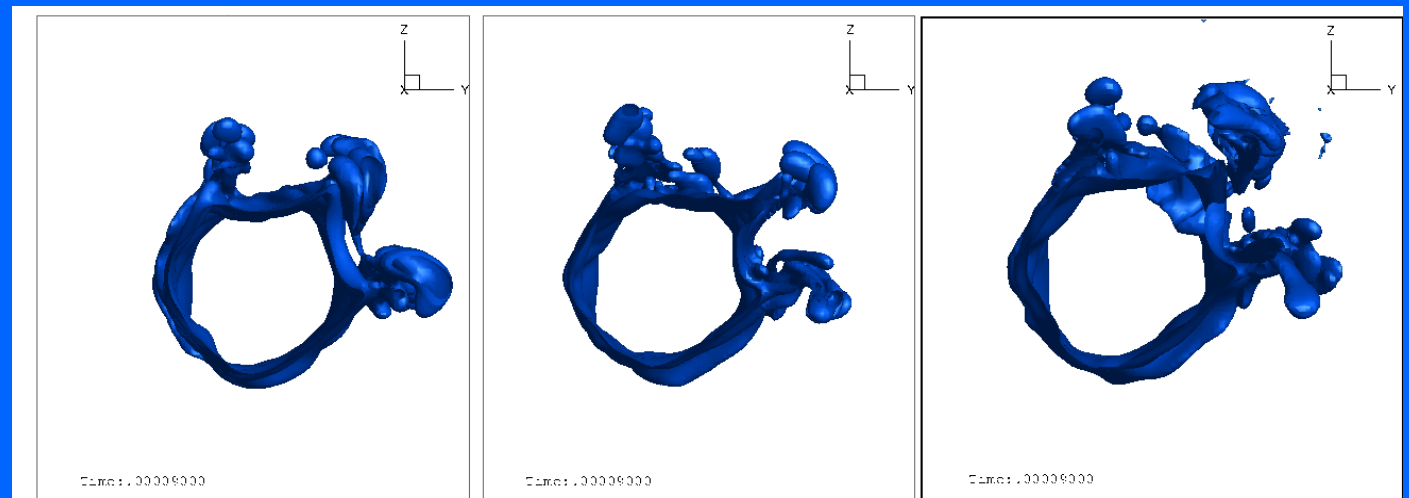
*Sin 3θ*

*Sin 5θ*

*Axisymmetric*

*Initial perturbation*

$t = 90 \mu s$





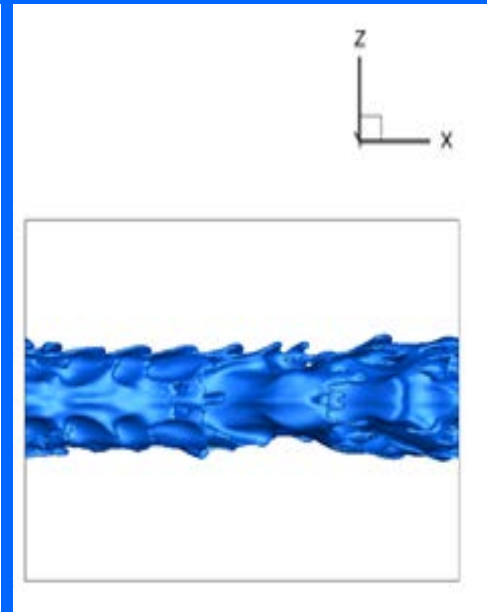
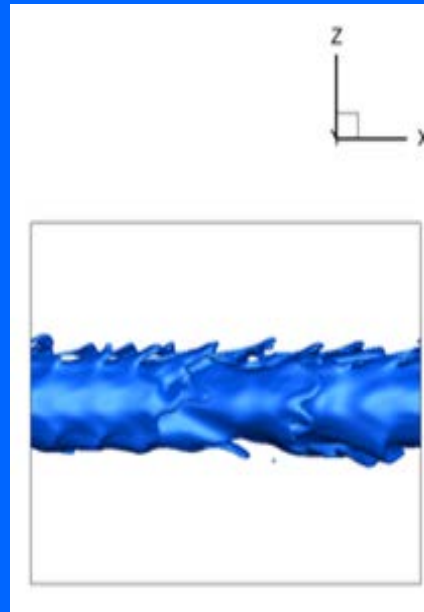
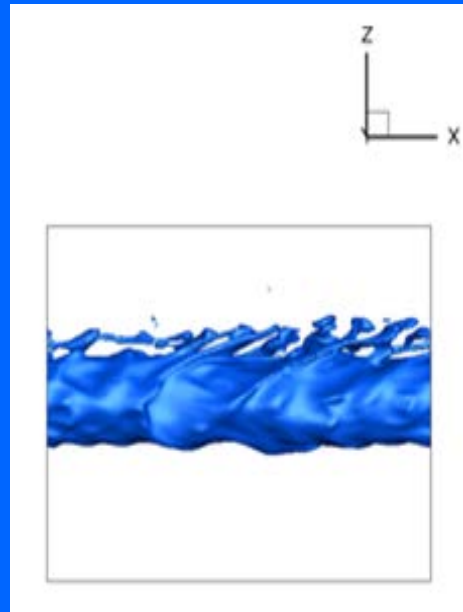
# Density Ratio affects 3D instability

$Re = 1600$

$We = 230,000$

$\mu_1 / \mu_2 = 9 \times 10^{-4}$

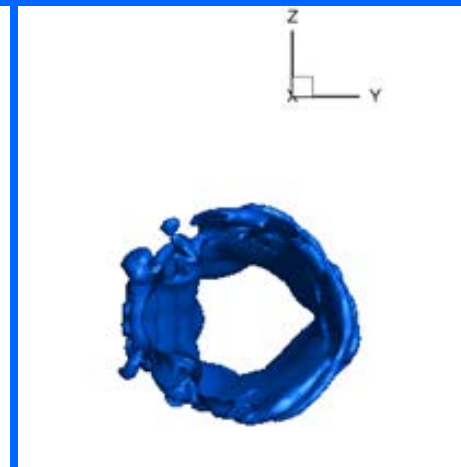
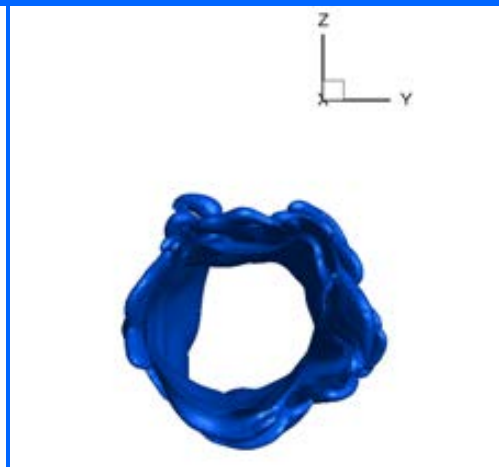
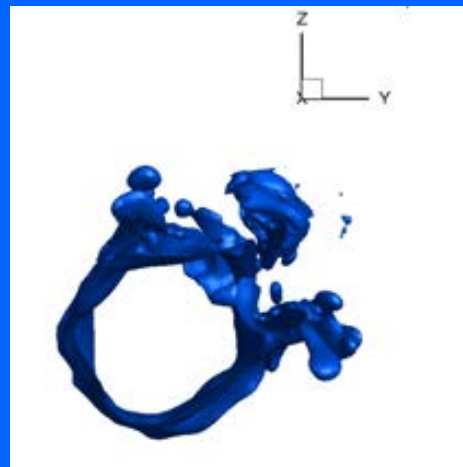
$t = 90 \mu s$



$\rho_1 / \rho_2 = 0.1$  ,

$0.5$  ,

$0.9$



# Vorticity

Originally axisymmetric; rings tilt; streamwise vorticity develops.

Vorticity magnitude increases with time.

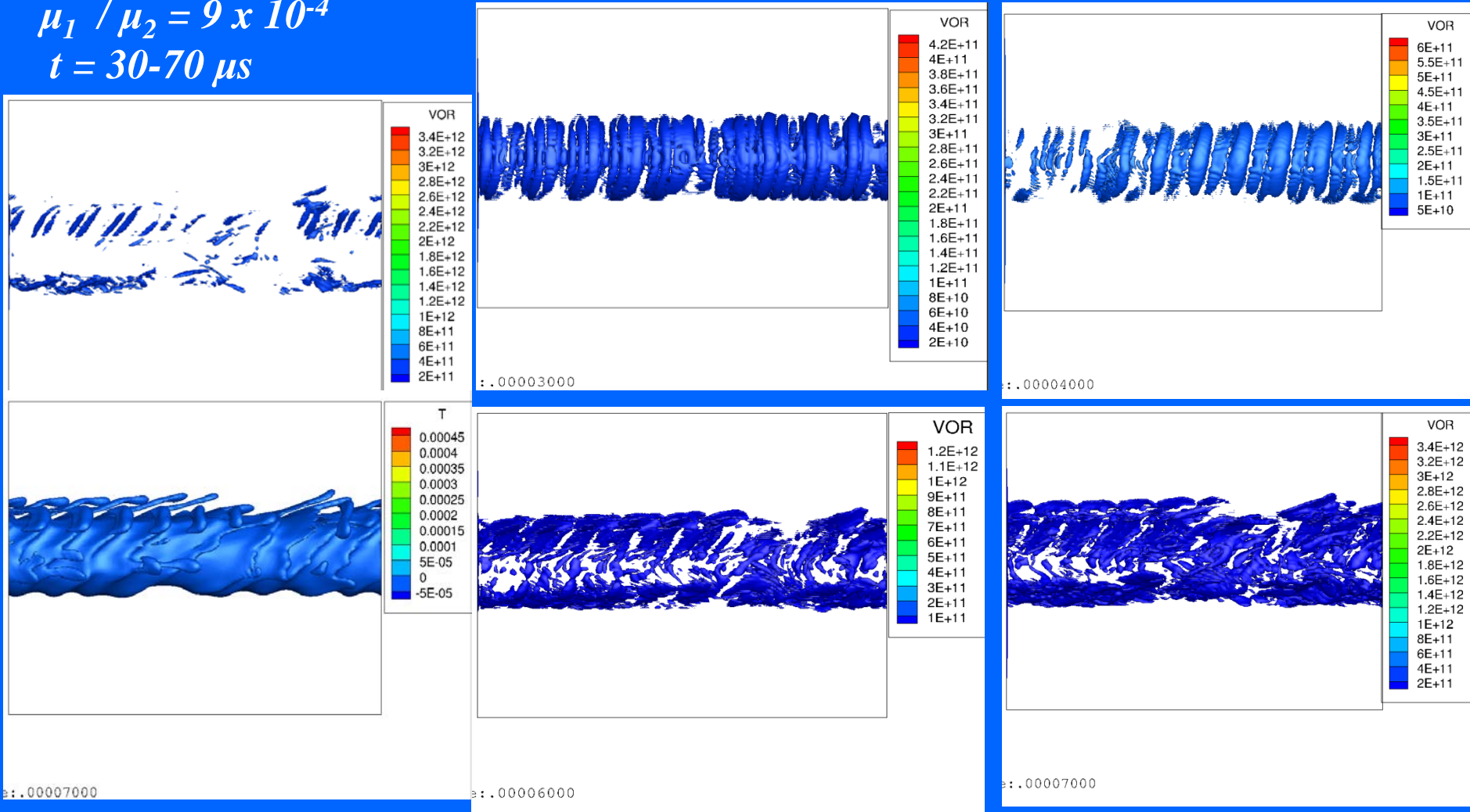
Maximum vorticity in braid region.

$$Re = 1600$$

$$We = 230,000$$

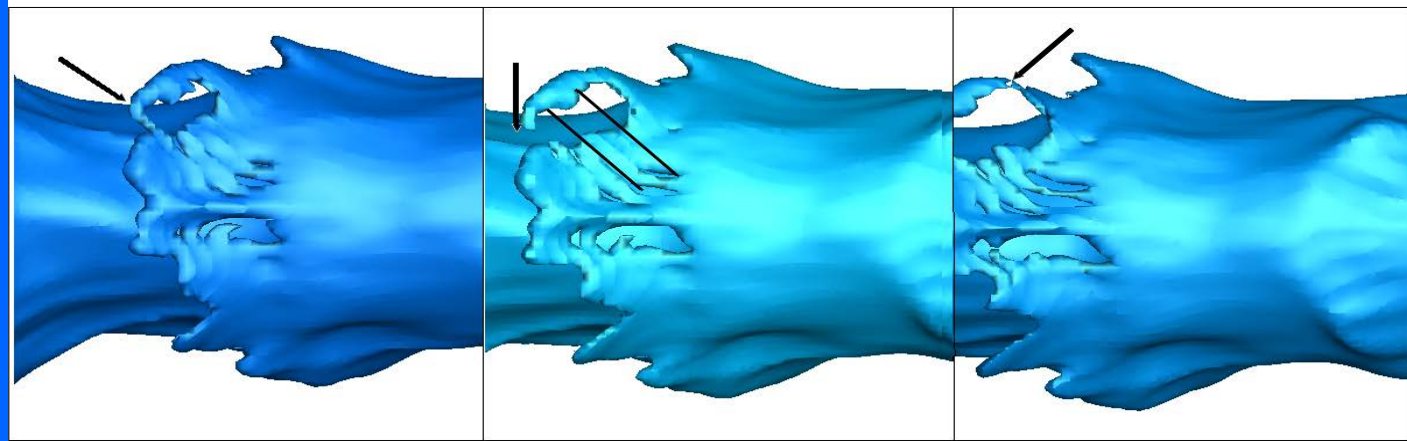
$$\mu_1 / \mu_2 = 9 \times 10^{-4}$$

$$t = 30-70 \mu s$$



# *Lobe and ligaments*

*time* →



- $We = 230000$

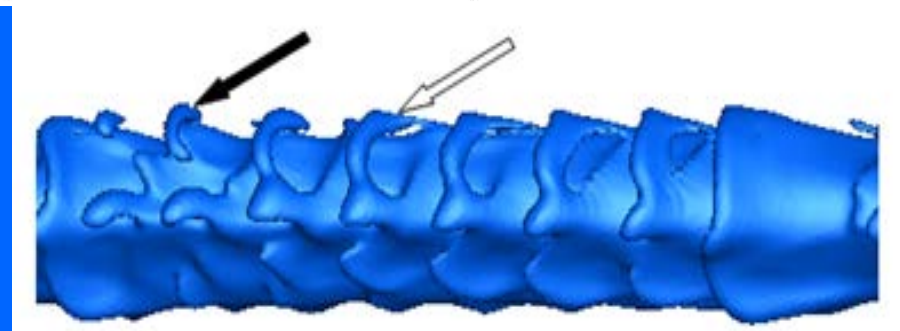
- $Re = 1600$

- $\rho_1 / \rho_2 = 0.1$

- $\mu_1 / \mu_2 = 9 \times 10^{-4}$

- $We = 23000, Re = 1600,$

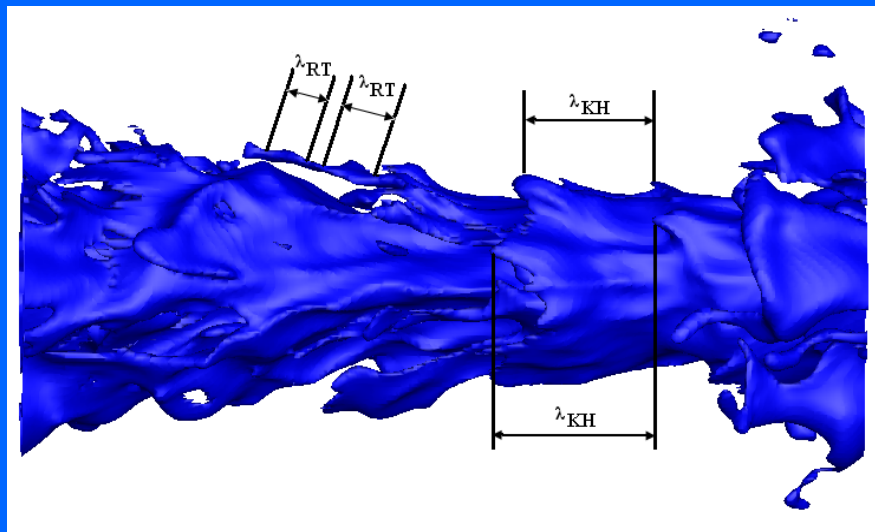
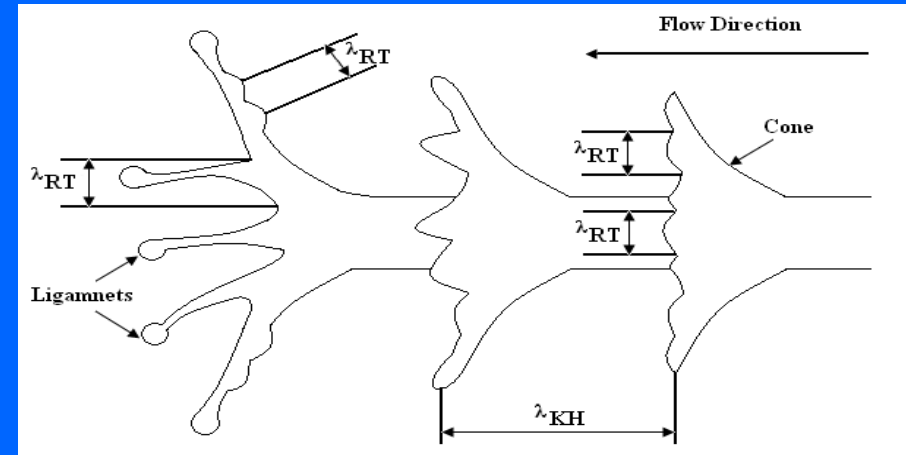
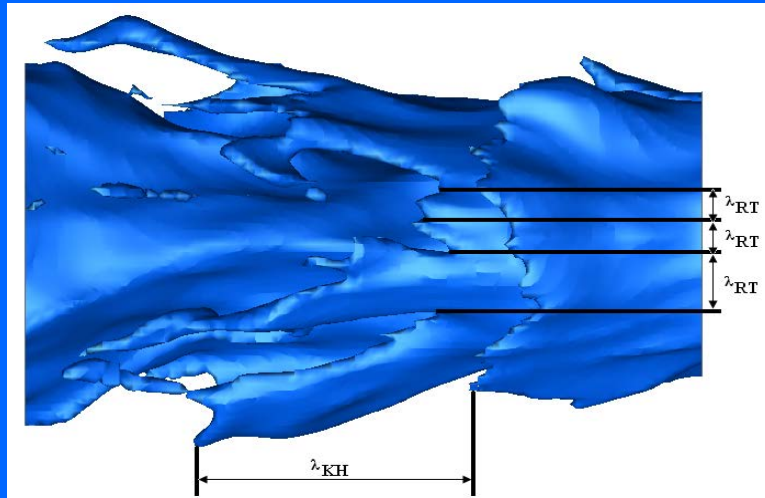
- $\rho_1 / \rho_2 = 0.1, \mu_1 / \mu_2 = 9 \times 10^{-4}$



- The cone-shaped liquid surface is sheared at the corners of the lobe where the maximum vortical motion occurs.
- The stretching thins the lobes yielding holes.
- The detached liquid from opened hole produces two ligaments that stretch away from core into the gas.
- The capillary effects can break the ligaments into droplets.

# *Wavelength Analysis on Unstable Liquid Jet*

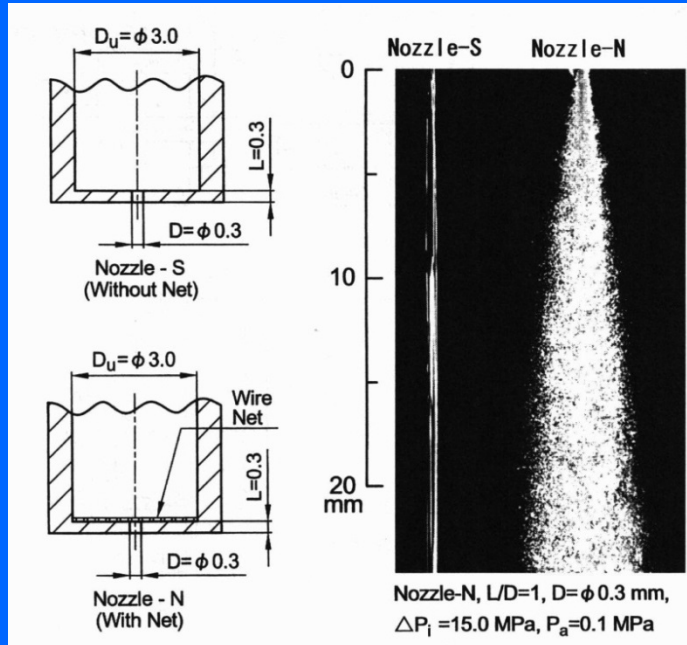
**K-H waves form cones, azimuthal R-T waves form on cones, ligaments result from R-T waves, capillary action plus more R-T waves on ligaments.**



- **Transverse R-T instability wavelengths on the crests varies between 20 to 70  $\mu\text{m}$ .**
- **R-T instability wavelengths on the ligaments vary between 15-50  $\mu\text{m}$ .**
- **Acceleration is the main factor to develop the R-T instability on the ligaments and crests .**

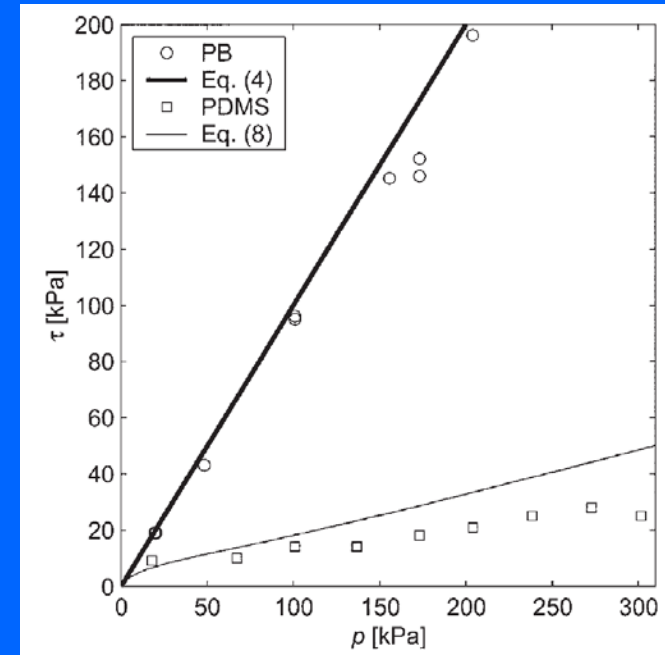
# *Cavitation in high pressure injectors*

**Cavitation occurs well above classical threshold pressure and has huge impact on jet breakup.**



← **Hiroyasu and co-workers (1998, 2000, 2001)**

**Kottke, Bair and Winer (2005)** →



**- - Total-stress criterion proposed by Winer and Bair (1987) and independently by Joseph (2001):**

$$T_{11} > -p_c$$

$$\mathbf{T} = \mu [(\nabla \mathbf{u}) + (\nabla \mathbf{u})^T] - p\mathbf{I}$$

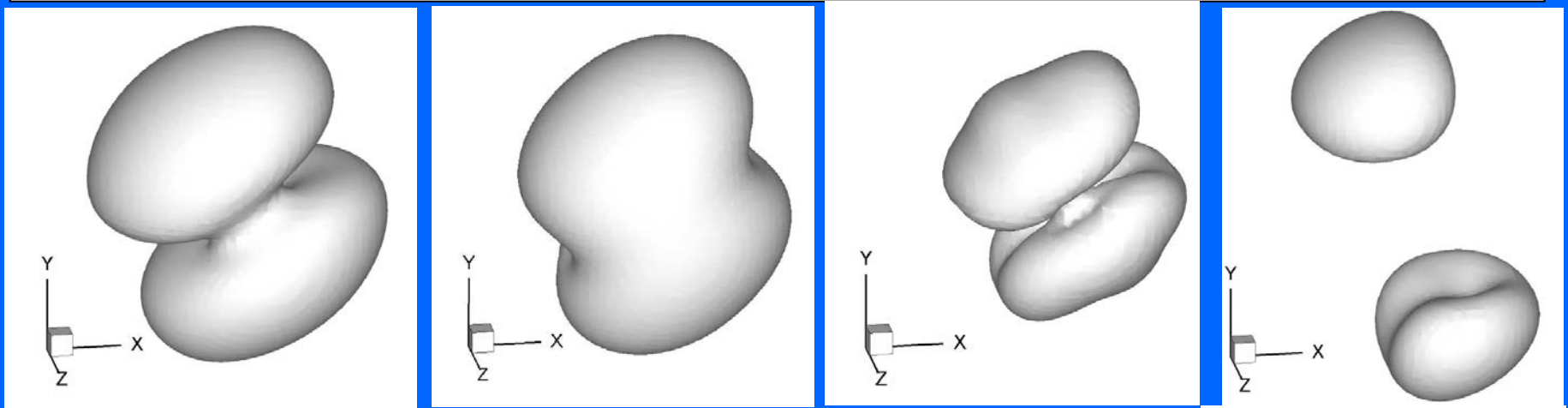
**(Maximum principal stress)**



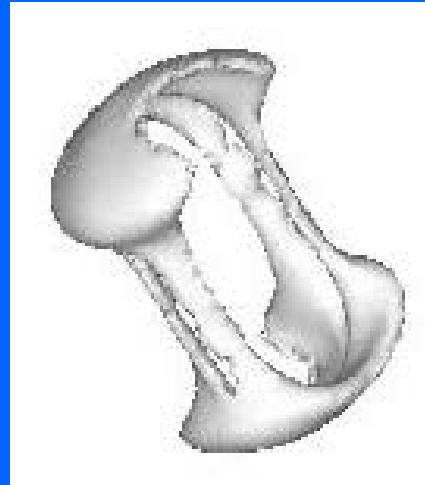
# *Calculation of bubble dynamics and far-field impact in strained flow*

**Dabiri, Sirignano, and Joseph (2007, 2008, 2010a, 2010b)- NS /LS**

**Bubble shape after rebound. From left to right: Case 1,  $Re=1.0$ ,  $We=0.1$ ; Case 2,  $Re=0.5$ ,  $We=0.1$ ; Case 3,  $Re=1.0$ ,  $We=\infty$ ; Case 4,  $Re=1.242$ ,  $We=0.0298$ .**



**In some cases with greater pressure variation, holes were formed upon collapse; threads developed connecting two masses and then broke.**



# *Summary*

**Stretching of liquid sheets can result in tearing and hole formation.**

**Vorticity dynamics explain liquid jet deformation. We can learn from experiences with single-phase jets and should seek to unify the single-phase and multi-phase phenomena.**

**Density differences and surface tension modify vorticity. R-T instability is best explained through vorticity dynamics.**

**Variation in pressure can cause changes in capillary instability and vorticity dynamics. Still, cavitation is possible at high injector pressures.**

*THANK YOU.*